

Plasma II L12: Plasma diagnostics 1

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- Introduction
- Passive measurements
 - EM emission
 - Bremsstrahlung
 - Line emission
 - Synchrotron radiation
 - Cyclotron emission
 - Magnetic fields
- Active measurements
 - Langmuir probes
 - Interferometry
 - Thomson scattering
- Appendix
- I.H. Hutchinson, “*Principles of plasma diagnostics*”, Second edition, Cambridge University Press



Plasma diagnostics 1

Plasma diagnostics 2

- Progress in physics usually through quantitative confrontation between theory and experiment
 - Requires theory predictions in realistic settings
 - Requires accurate and extensive measurements of the (plasma) properties → **Diagnostics**

- Ions (various species and charge states), electrons and neutrals

$f_s(\bar{x}, \bar{v}, t)$ - *Distribution function for particle species s*

- Moments of the distribution function (see L2)
 - $N=0$: Density
 - $N=1$: Flow velocity
 - $N=2$: Pressure/Temperature
- Deviations from a Maxwell-Boltzmann distribution (e.g. fusion alphas!)
- Impurity concentrations
- Magnetic and electric fields
 - Their sources: electric currents and electric charges/potential
- Time-averaged values or fluctuations

Approaches to diagnostics

- What do we have to measure to deduce a particular quantity?
- What can we conclude from a specific measurement?

- Physics
- Detection
 - Scheme
 - Technology
 - Perturbation of the plasma
- Signal-treatment/transfer
- Signal acquisition
- Post processing
- Data storage



Address only these aspects in this course

Passive measurements

- Radiation diagnostics
 - Spectroscopy
 - Bolometry
 - Radiation in magnetised plasmas
- Magnetic measurements
- Ion loss detectors
- ...

Active measurements

- Langmuir probes
- Interferometry
- Scattering diagnostics
- Neutral beam (e.g. CXRS)
- Heavy beam probes
- ...

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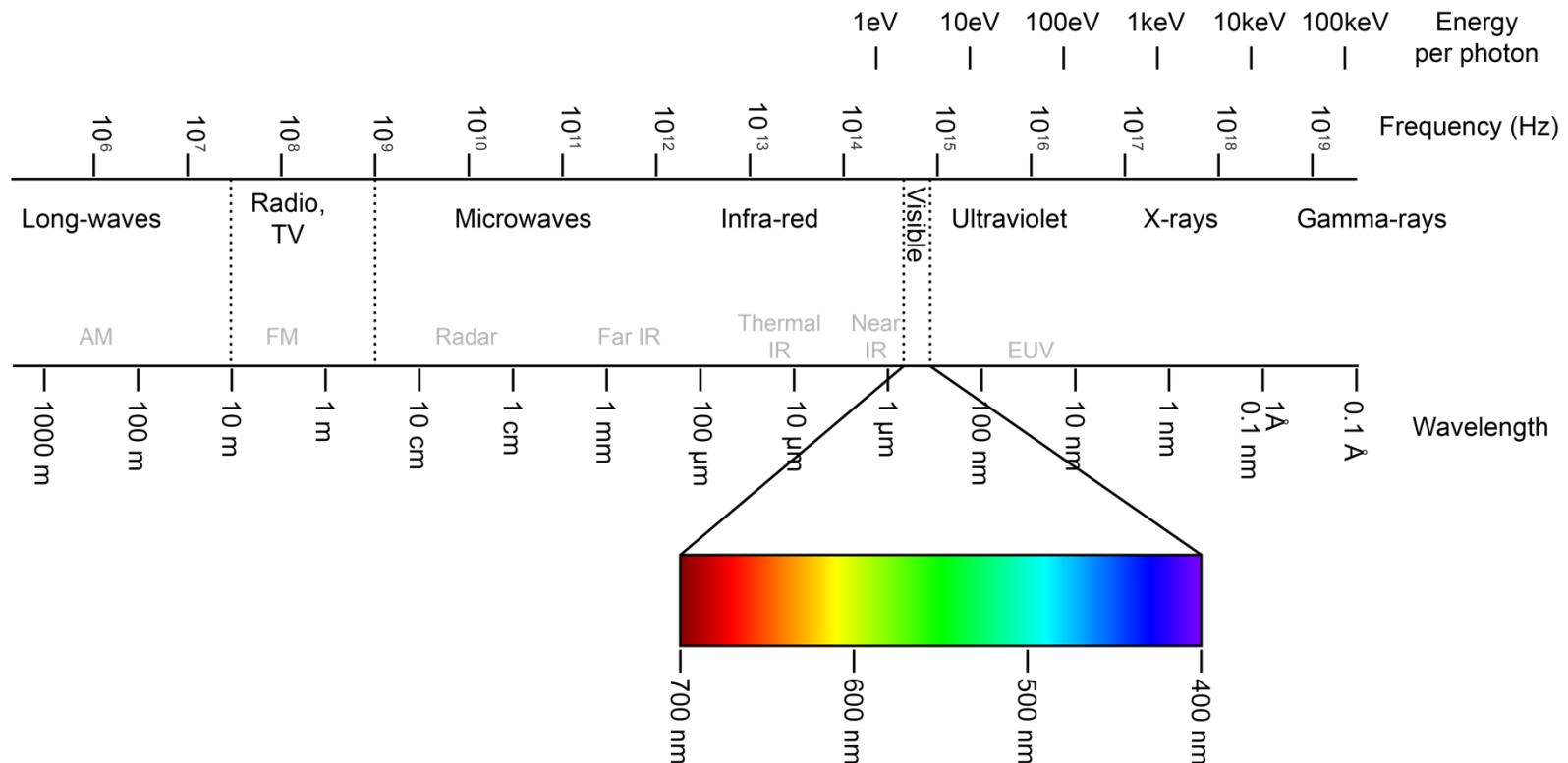


Plasma diagnostics 1



Plasma diagnostics 2

EM spectrum



- Plasmas can emit electro-magnetic radiation over a wide range of frequencies/wavelengths
 - From hard X-rays to cm waves
 - Emission dominated by characteristic lines on top of a lower continuum
 - (Laboratory) Plasmas often/for most ν optically thin → Measure $S_\nu = \int \epsilon_\nu(\mathbf{r}) d\mathbf{r}$
- Electrons
 - Free-free collisions with ions & electrons → Bremsstrahlung (continuum)
 - Free-bound → Recombination (continuum)
 - Bound-bound transitions → Line radiation
 - Free in B-field → Synchrotron & cyclotron radiation

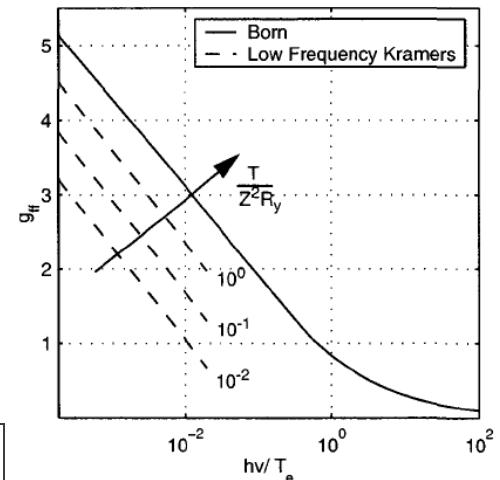
- Coulomb collisions deflect, i.e. accelerate charges → e-m emission

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$$j(\nu) = \frac{n_e n_i Z^2}{4\pi} \left(\frac{e^2}{4\pi\epsilon_0} \right)^3 \frac{32\pi^2}{3\sqrt{3}m_e^2 c^3} \left(\frac{2m_e}{\pi T_e} \right)^{1/2} e^{-h\nu/T_e} \bar{g}$$

- With Maxwell-averaged Gaunt factor \bar{g}
- For $0.1 \leq h\nu/T_e \leq 10$, \bar{g} of order unity → temperature dependence of $j(\nu)$ well described by $\propto T_e^{-1/2} e^{-h\nu/T_e}$

Total emissivity



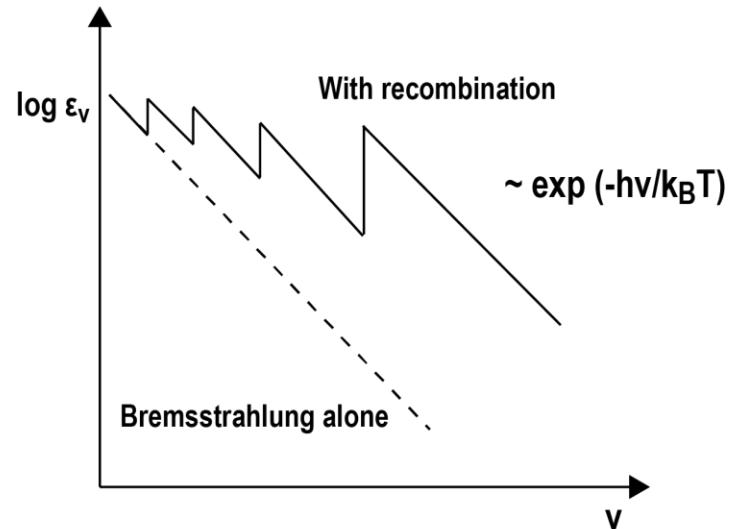
- Intensity and spectrum of bremsstrahlung yields information on plasma density, charge state and temperature

Recombination radiation

- Energy balance (hydrogen)

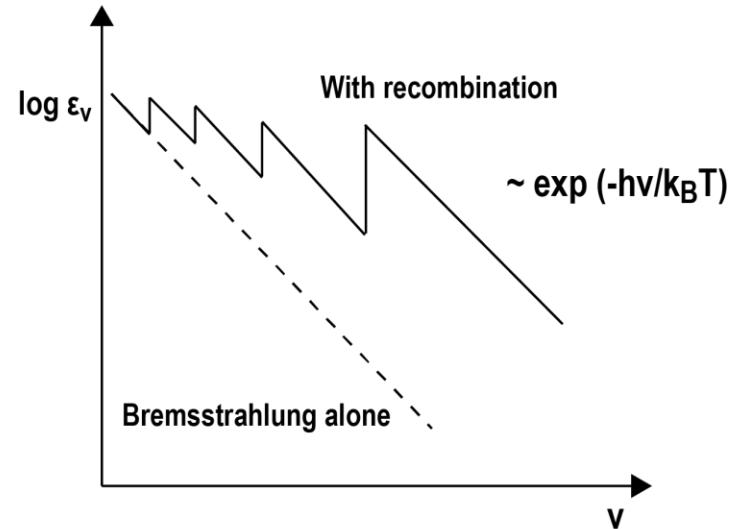
$$h\nu = \frac{1}{2}mv^2 + \frac{R_y}{n^2}$$

- Introduces 'edges' in the continuum as final states with lower n increase photon energy in discrete steps
 - For $h\nu \ll R_y$ the contribution of recombination is negligible



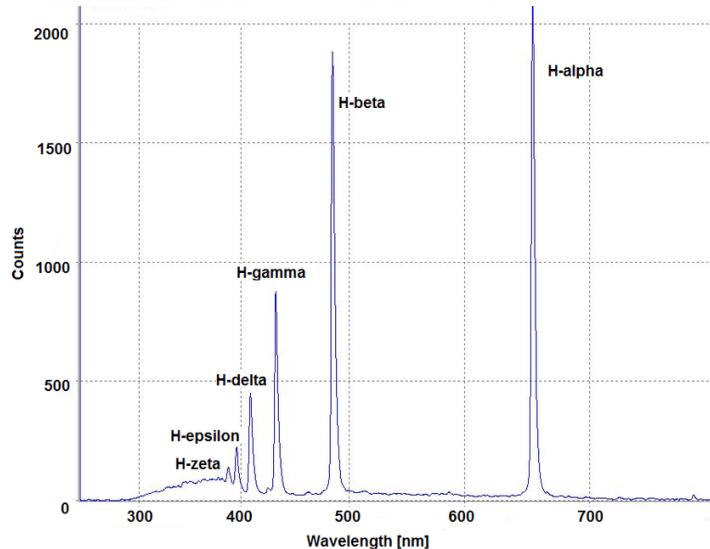
Use Bremsstrahlung to determine electron temperature

- Compare signal intensity with different filters → “Filter temperature”
 - Detect in the soft-x-ray regime
 - Use Be filters to suppress ‘low’ frequencies



- Pattern for hydrogen recognised by J. Rydberg in 1888
 - Express emission through wavenumber λ^{-1} was key for discovery

Visible emission of a hydrogen plasma in a theta pinch



For hydrogen

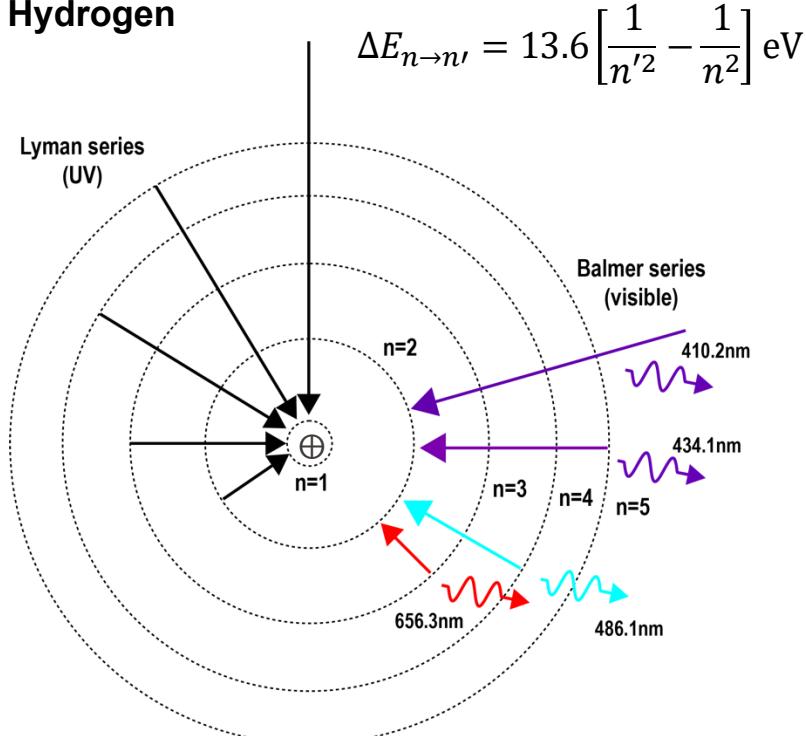
$$\frac{1}{\lambda} = R_H \left(\frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

Line-radiation: Frequency selection

- Transitions between bound electron states lead to discrete photon energies

$$\Delta E_{n \rightarrow n'} = h\nu = \frac{hc}{\lambda} = E_n - E_{n'}$$

Ex.: Hydrogen



Line-radiation: Line intensity

- Emission depends on density n_n of bound electrons in state n and transition probability $A_{n \rightarrow n'}$

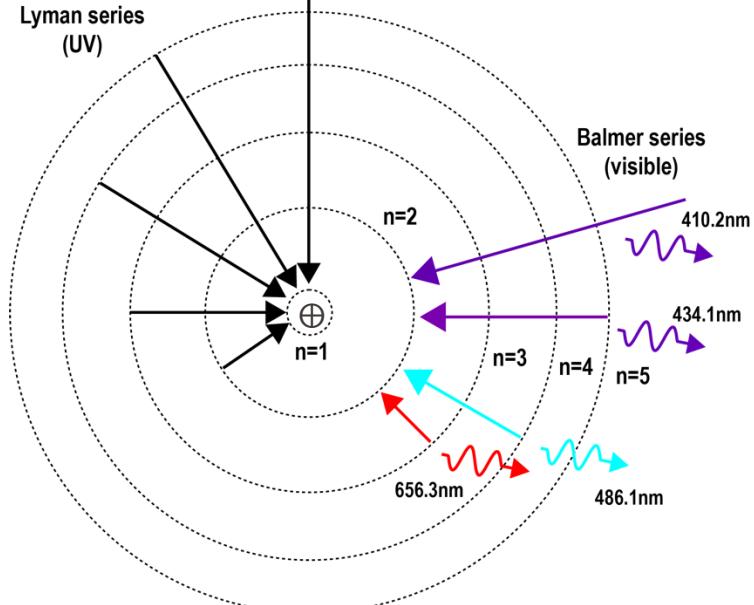
$$\varepsilon_L = \Delta E_{n \rightarrow n'} A_{n \rightarrow n'} n_n$$

- Measured line intensity (radiance)

$$I_L = \frac{1}{4\pi} \int \varepsilon_L dl$$

Ex.: Hydrogen

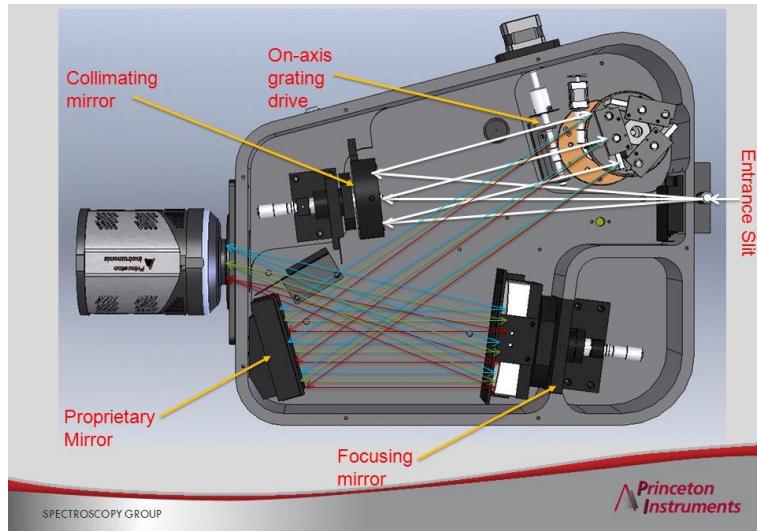
$$\Delta E_{n \rightarrow n'} = 13.6 \left[\frac{1}{n'^2} - \frac{1}{n^2} \right] \text{ eV}$$



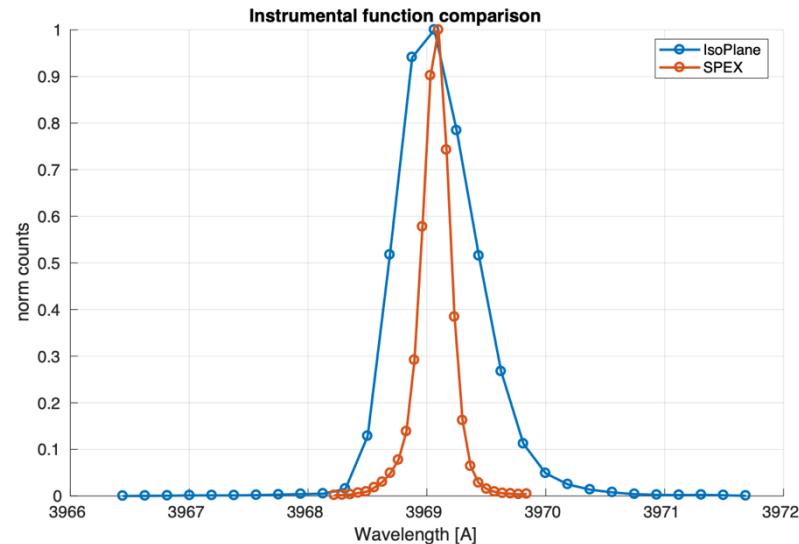
- Use a grating to ‘Fourier transform’ the emitted light

Spectrometer

- Schmidt-Czerny-Turner Spectrograph



- Ex.: Instrumental function of TCV's divertor spectrometers

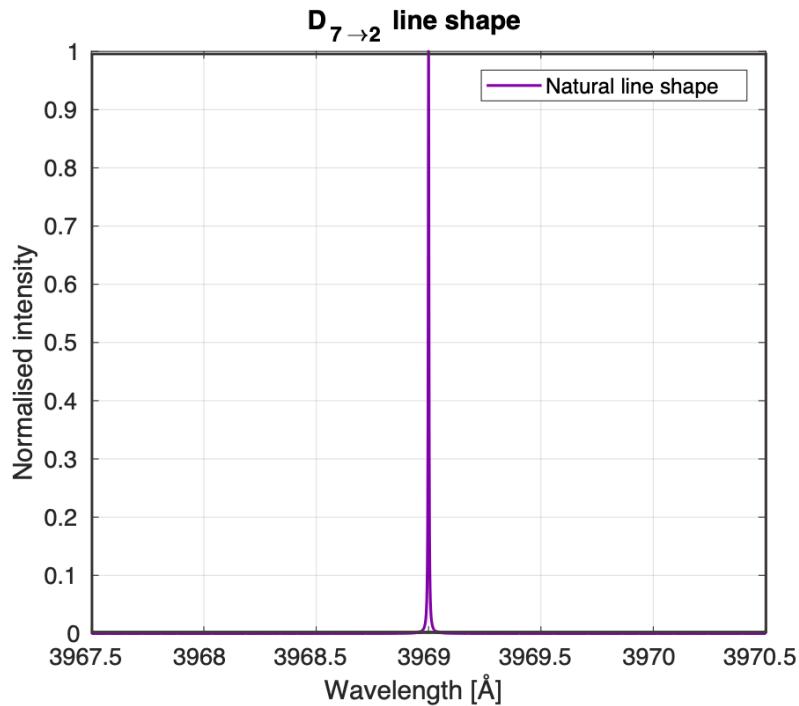


[Courtesy of L. Martinelli]

Spectral line shape

- Natural broadening causes a *Lorentzian distribution* of the emission intensity
 - Width inversely proportional to transition probability (Einstein coefficients)
- **Ex.:** D- ϵ

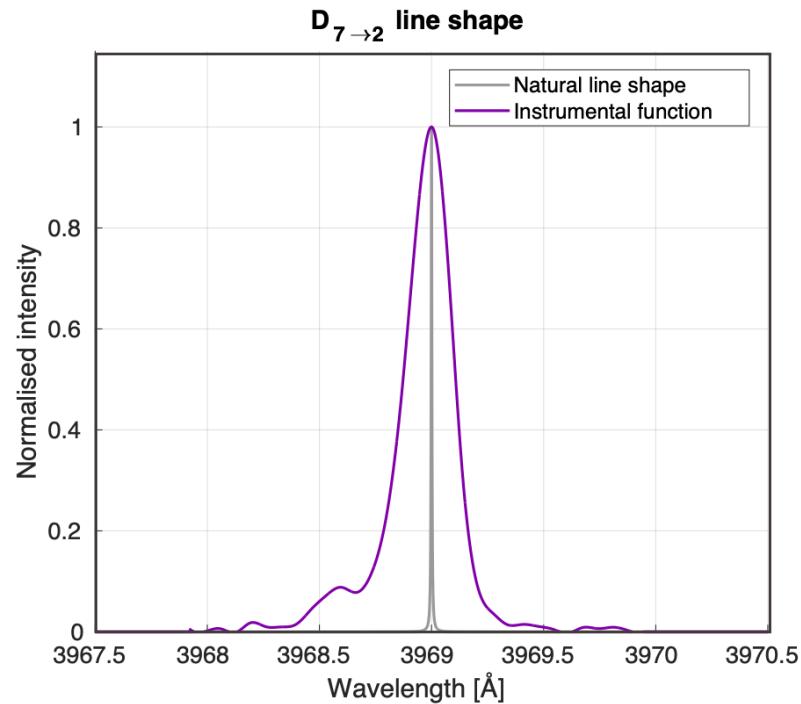
$$\text{FWHM}(D_{7 \rightarrow 2}) \simeq 4 \cdot 10^{-3} \text{\AA}$$



Spectral line shape

- Diagnostic setup broadens the line (*instrumental broadening*)
- **Ex.:** TCV high resolution divertor spectrometer (SPEX)

$$\text{FWHM}(\text{SPEX}) \simeq 0.20\text{\AA}$$



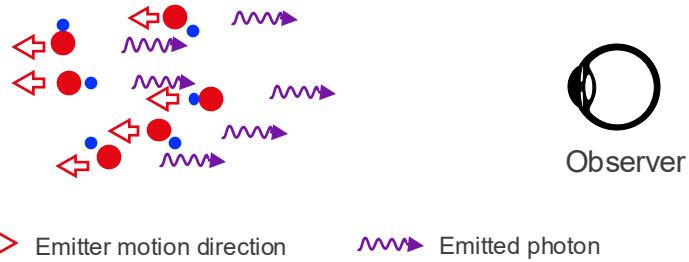
[Courtesy of L. Martinelli]

Spectral line shape

... can reveal the bulk motion

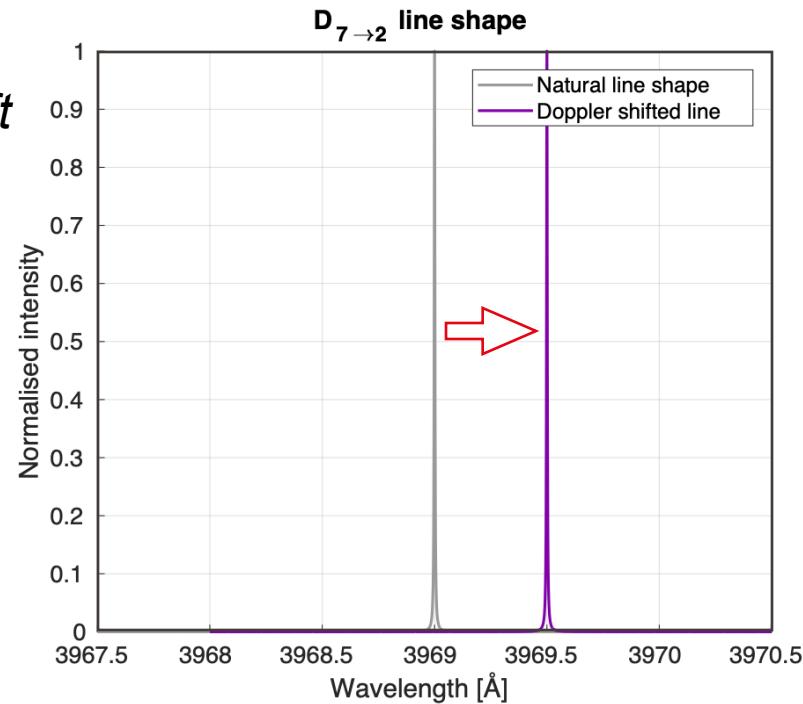
- Bulk motion of emitting species along the line of sight causes a *Doppler shift*

$$\Delta\lambda = \lambda_0 \frac{v}{c}$$



- Ex.: TCV plasma moving with ~35km/s with respect to spectrometer

$$\Delta\lambda_{D-\varepsilon} \simeq 0.5\text{\AA}$$

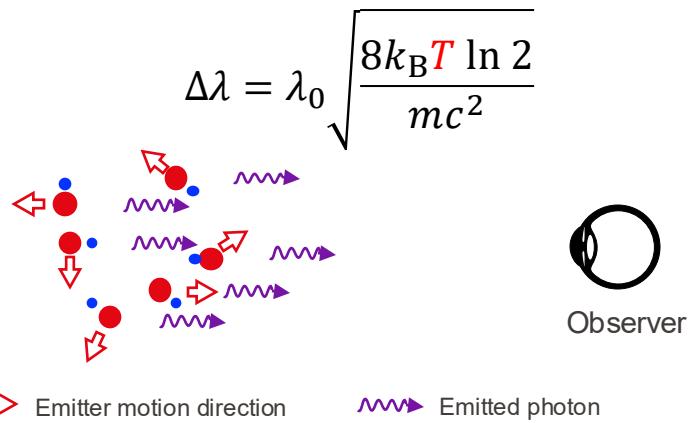


[Courtesy of L. Martinelli]

Spectral line shape

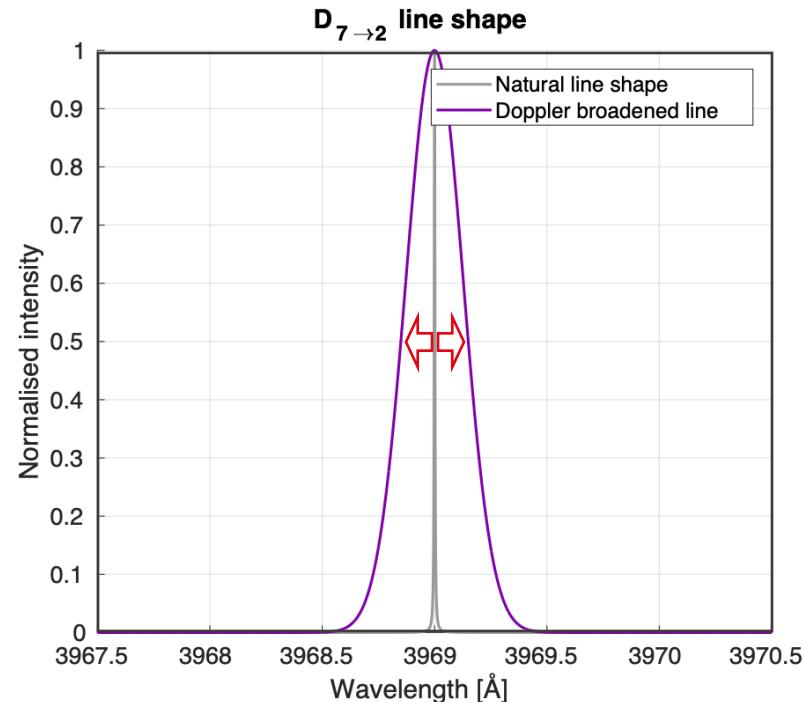
... can reveal the atom temperature

- Thermal motion of emitting particles leads to *Doppler* broadening



- Ex.: 1eV deuterium plasma

$$\Delta\lambda_{\text{FWHM}} \simeq 0.25 \text{ \AA}$$

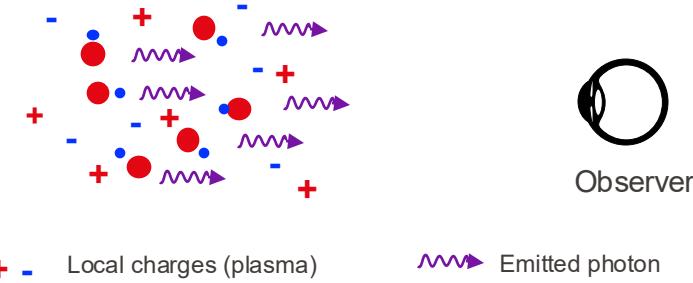


[Courtesy of L. Martinelli]

Spectral line shape

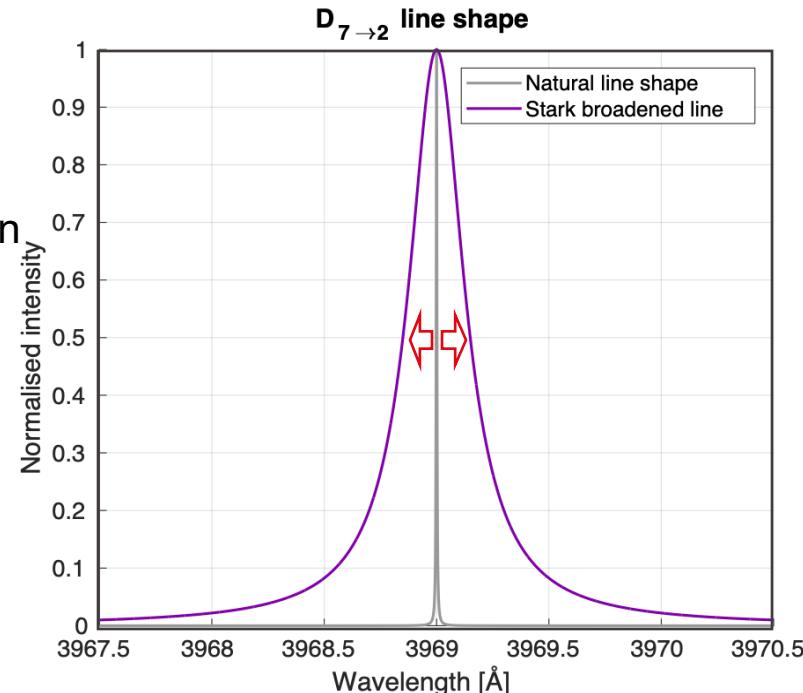
... can reveal the plasma density

- Random electric microfields induced by surrounding plasma lead to *Stark broadening*
 - Approximated by *Lorentzian* distribution



■ Ex.: D- ϵ

$$\text{FWHM}_{\text{Stark}}(D_{7 \rightarrow 2}) \approx 0.21 \cdot \left(\frac{n_e [m^{-3}]}{10^{19}} \right)^{2/3} \text{\AA}$$



[Courtesy of L. Martinelli]

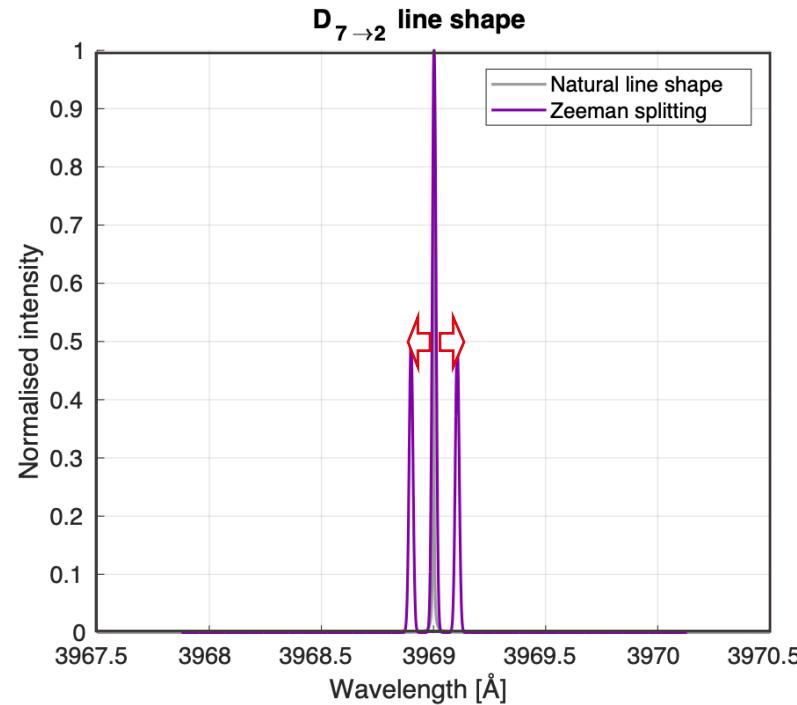
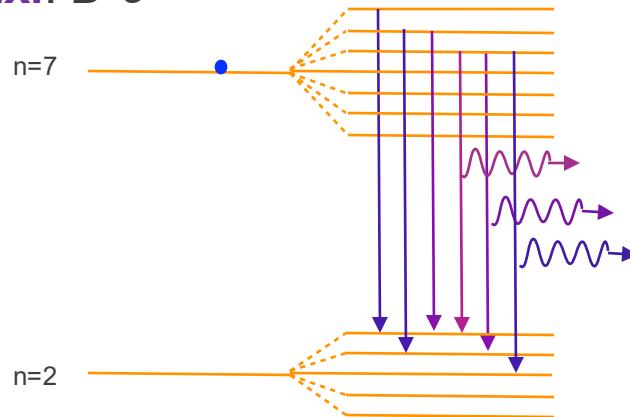
Spectral line shape

... can reveal the magnetic field

- Magnetic field perpendicular to the LOS splits upper and lower energy levels (Zeeman splitting)

$$\Delta E = \mu_B (L_Z + S_Z) \cdot \mathbf{B} , \quad \mu_B = \frac{\hbar e}{2m_e}$$

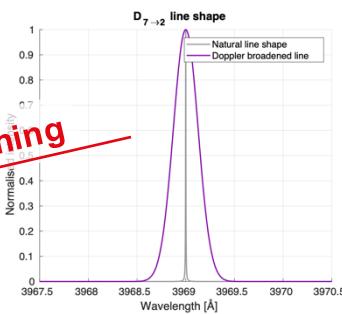
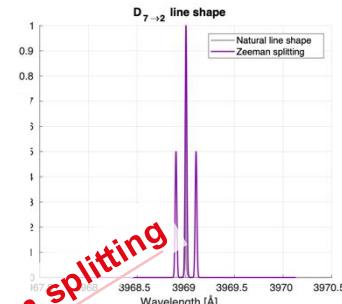
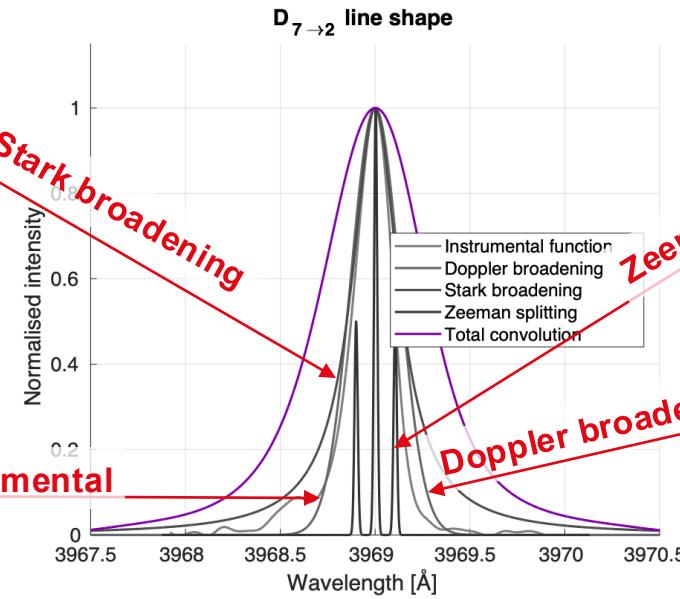
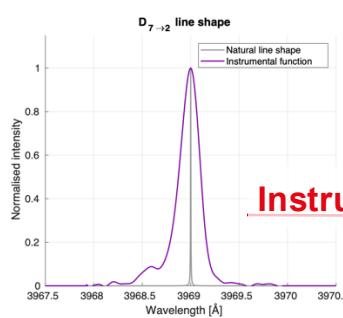
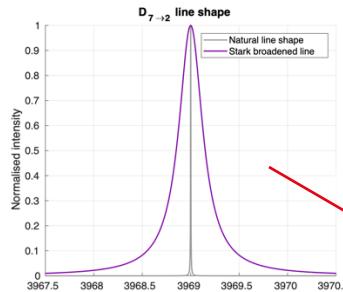
- Ex.: D- ϵ



[Courtesy of L. Martinelli]

Measured spectral line shape is a convolution of all effects

- Knowledge of instrumental function, density, temperature and/or magnetic field help to constrain fit



Spectroscopy

Hydrogenic lines of interest

- Balmer lines ($n' = 2$) optically thin
- Lyman lines ($n' = 1$) radiate more power, but are all in the UV
 - Many lens and window materials do not transmit in the UV

Balmer lines	Wavelength [nm]	Energy difference [eV]	Colour
Alpha (3→2)	656.3	1.89	red
Beta (4→2)	486.1	2.55	aqua
Gamma (5→2)	434.0	2.86	blue
Delta (6→2)	410.2	3.03	violet
Epsilon (7→2)	397.0	3.13	(ultraviolet)
Zeta (8→2)	388.9	3.19	
Eta (9→2)	383.5	3.23	
Balmer break (Inf→2)	364.6	3.40	

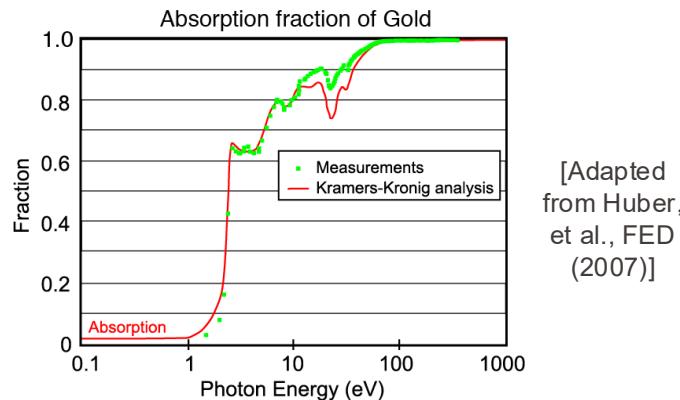
Line intensities can be used to infer densities and temperatures

- Collisional-radiative modelling
 - Excited states n created by collisions of electrons with lower states or through recombination of ions and electrons
 - Example: Balmer lines

$$\varepsilon_{n \rightarrow 2} = n_e n_{H^0} PEC_{n \rightarrow 2}^{\text{EX}}(n_e, T_e) + n_e n_{H^+} PEC_{n \rightarrow 2}^{\text{RC}}(n_e, T_e)$$

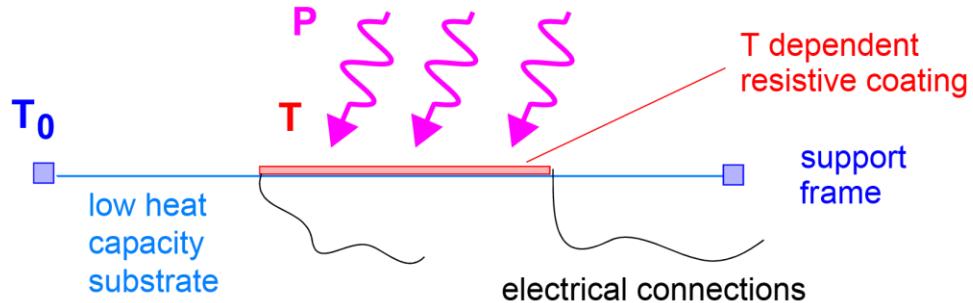
- Photo emission coefficients (PEC) depend on cross-sections and transition probabilities → tabulated in databases (e.g. ADAS)
- Use 3 Balmer line measurements to solve for the unknowns $n_e \sim n_{H^+}, T_e, n_{H^0}$

- Measurement of **radiated power**, regardless of wavelength, referred to as ***bolometry***
 - Eg. for power balance
- Requires a detector that is sensitive to a broad spectrum (in MCF typically from soft X-ray to the visible)
 - Metal foil (can be 'blackened')



[Sheikh, et al., RSI (2016)]

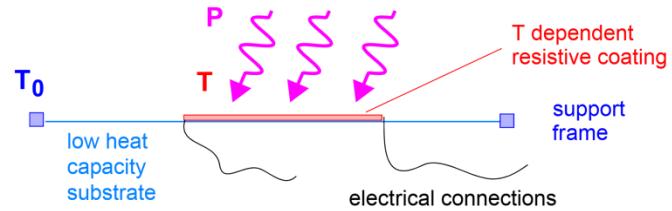
- Temperature increase of metal foil yields radiance



[Adapted from H. Weisen]

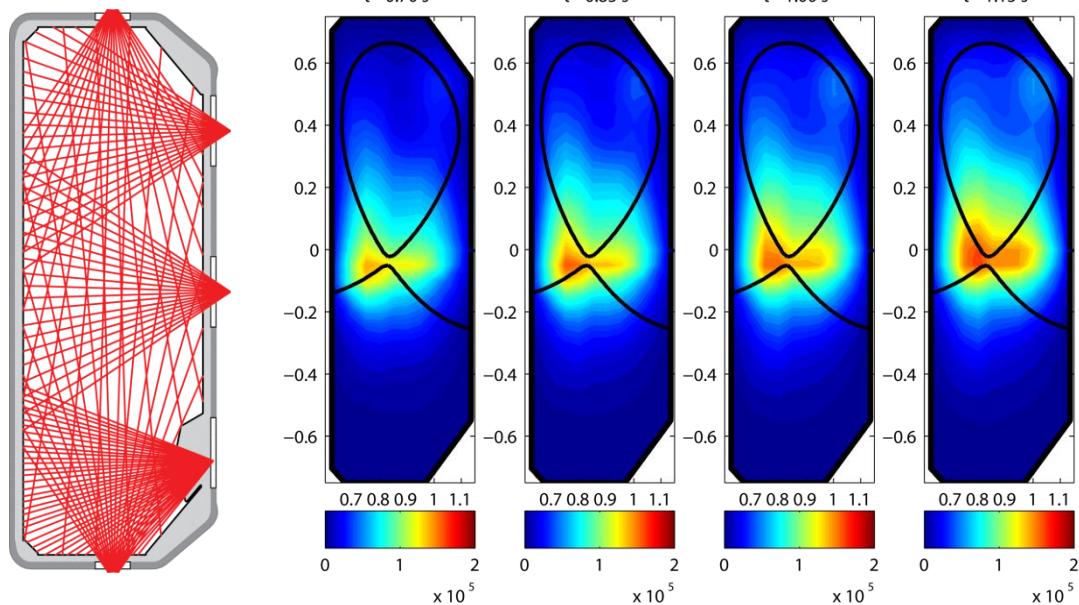
- Use temperature dependence of electrical resistivity
 - Heat capacity and thermal contact to heat sink limits time response

- Sensitivity and time response



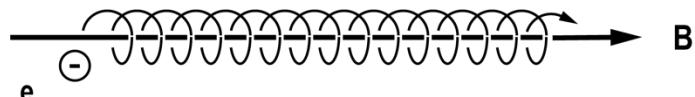
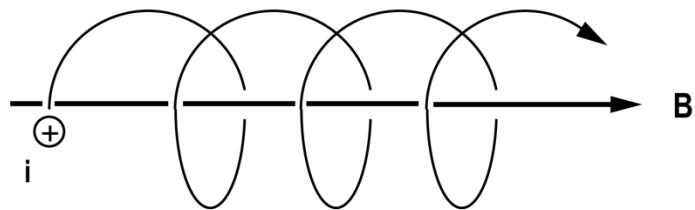
Inversion of line integrated measurements

- Emission in an optical thin medium leads to line-integrated measurements → need an inversion (tomography) for spatial resolution
 - Use intersecting lines-of-sight (LoS) and/or constrain geometry of solution
- **Ex.:** TCV bolometer
 - 120 LoS



Additional radiation in magnetised plasmas

- Charged particles that move perpendicular to a magnetic field are subject to the Lorentz force and, hence, an acceleration and emit even without collisions (*magneto-bremsstrahlung*) → **Cyclotron radiation**
 - When particles move with relativistic speeds *cyclotron radiation* is also referred to as *synchrotron radiation*



- Main contribution from electrons

$$\omega_{c,e} = l \cdot \frac{eB}{m_e} \quad \text{with } l = 1, 2, 3, \dots$$

- Cyclotron radiation is optically thick → Black-body radiation → Intensity depends only on electron temperature!
- With $\hbar\omega_{c,e} \ll k_B T_e$ the ‘black-body radiation’ is described by Rayleigh-Jeans law

$$I_\nu^{\text{bb}} = \frac{\omega_{c,e}^2 k_B T_e}{8\pi^3 c^2}$$

Electron-cyclotron emission spectrum

- $\omega_{c,e}$ depends on magnetic field strength and, hence in tokamaks on the major radius
 - Emission at $\omega_{c,e}$ characteristic for a location → measurements at several frequencies yield T_e -profile

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Plasma diagnostics 1



Plasma diagnostics 2

- Magnitude of magnetic fields first measured by Carl Friedrich Gauss in 1833

Examples of B-field diagnostics

- Pick-up coils & flux loops (magnetic induction)
- Hall probe
- Fluxgate-magnetometer
- Spectroscopy → see Zeeman splitting

Measure B-field using magnetic induction

- Faraday's law $\nabla \times \mathbf{E} = -\frac{d\mathbf{B}}{dt}$

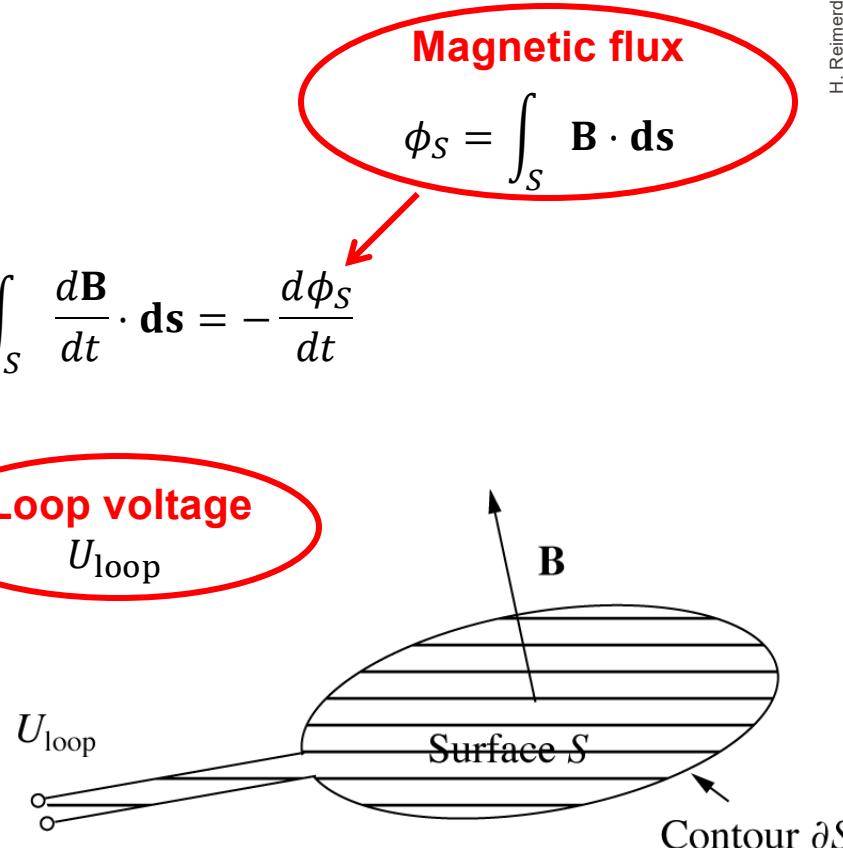
- Integral form

$$\oint_{\partial S} \mathbf{E} \cdot d\mathbf{l} = - \int_S \frac{d\mathbf{B}}{dt} \cdot d\mathbf{s} = -\frac{d\phi_S}{dt}$$

- Voltage induced in a loop (contouring the surface)

$$U_{\text{loop}} = -\frac{d\phi_S}{dt}$$

measures the component of $d\mathbf{B}/dt$ normal to (and integrated over) the plane of the loop



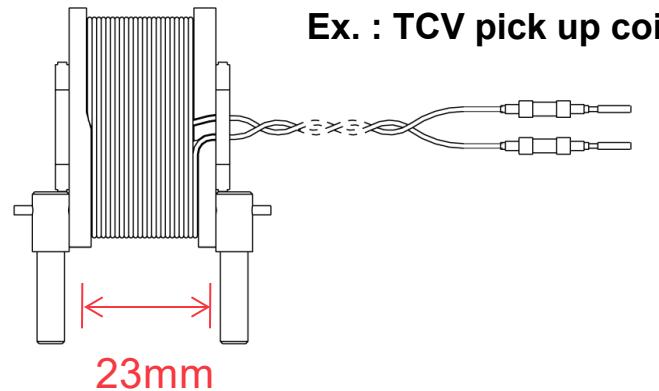
Various geometries exist

- For a small rigid probe ($|\mathbf{B}|/(\nabla \mathbf{B})_{\perp} \gg L_{\text{probe}}$) with a cross-sectional area A and N turns

$$U_{\text{probe}} = -NA \frac{dB_{\perp}}{dt}$$

- Probe measures the magnetic field component perpendicular to the probe surface B_{\perp}

- Other geometries
 - **Flux loops:** single loops that span a large area



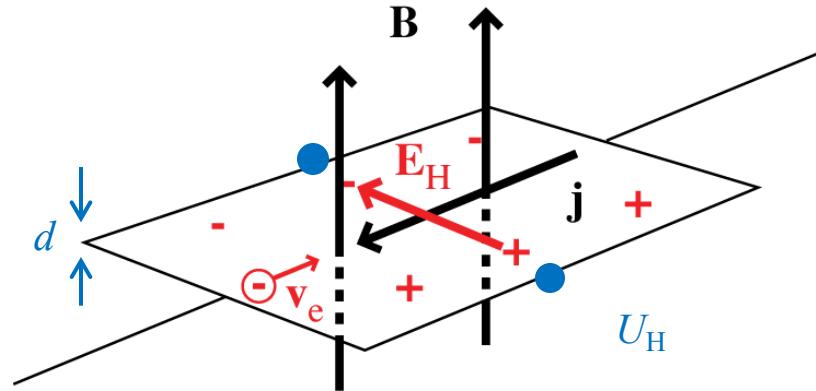
Measure B-field using the Hall effect

- Magnetic field leads to charge separation in a current carrying conductor

$$\mathbf{E}_H = -\mathbf{v} \times \mathbf{B} = -\frac{\mathbf{j} \times \mathbf{B}}{nq}$$

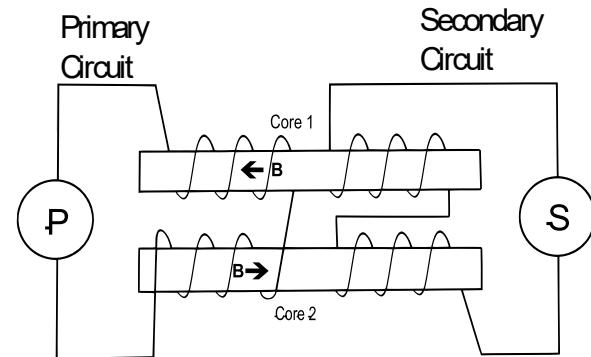
- Hall field depends on the the charge carrier density n and their sign q

- Advantages
 - Compact – miniaturization through IC technology



Measure B-field using Fluxgate-magnetometer

- Detect asymmetries in the saturation of a ferromagnetic coil core
- Constraints
 - Field of coil must be larger than measured field (not applicable to MCF)

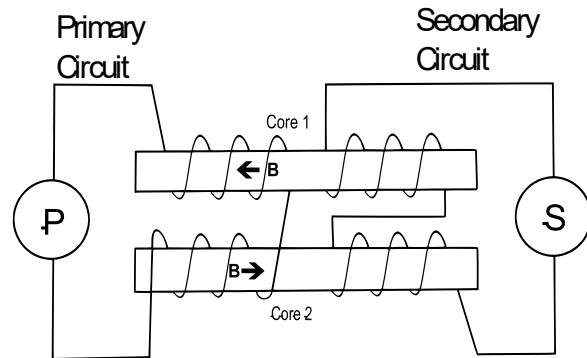


Source: Von Pdornel - Eigenes Werk, CC BY-SA 4.0,
<https://commons.wikimedia.org/w/index.php?curid=40110734>

- Without a magnetic field the voltages induced in the secondary loop (S) cancel

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 - Low power consumption
 - Rugged
 - Compact & light
 - Miniaturization recently advancing to the point of complete sensor solutions as IC



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Source: ESA

FGM diagnostic of Cluster satellites (see L2)

Textbooks

- J. Wesson, “Tokamaks”, Third edition, Oxford Science Publications – Chapter 10 ← *Good overview*
- I.H. Hutchinson, “Principles of plasma diagnostics”, Second edition, Cambridge University Press ← *Very detailed description based on first principles, goes well beyond course content*